

- Why Symmetries Relate Scattering/Decay Amplitudes

- Time Evolution: One solution of Schrödinger equation for time dependence of a state is $|\psi(t)\rangle = e^{-iHt} |\psi(0)\rangle$

where $H = \text{Hamiltonian}$

- Suppose there is an $SU(2)$ symmetry \vec{J}
 - + This could be spin, orbital angular momentum, isospin, etc
 - + Total \vec{J} operator is J^2 w/ eigenvalue $j(j+1)$
 - + 3rd component operator is J_3 w/ eigenvalue m
 - + Because this is a symmetry, an eigenstate $|j, m\rangle$ can only evolve to itself, i.e.,
$$e^{-iHt} |j, m\rangle = M_j |j, m\rangle$$

where $M_j = \text{constant}$

- + Also by the symmetry, M_j depends only on the quantum number j and not quantum number m . For example, if \vec{J} represents angular momentum, the symmetry is rotations. Spherical/rotational symmetry means we can rotate around the x axis to change where the z -axis points. That changes the m quantum number but can't otherwise change the time evolution.

- Transition amplitudes

- + The probability that we start w/ state $|\psi\rangle$ and measure state $|\phi\rangle$ at time t is

$$P = |\langle \phi | e^{-iHt} | \psi \rangle|^2$$

- + This means the probability amplitude is $\langle \phi | e^{-iHt} | \psi \rangle$

- + In a case with the $SU(2)$ symmetry, write each state as a superposition of $|j, m\rangle$ states

$$|\psi\rangle = a_1 |j_1, m_1\rangle + a_2 |j_2, m_2\rangle + \dots \quad |a_1|^2 + |a_2|^2 + \dots = 1$$

$$|\phi\rangle = b_1 |j_1, m_1\rangle + b_2 |j_2, m_2\rangle + \dots \quad |b_1|^2 + |b_2|^2 + \dots = 1$$

+ We know eigenstates w/ different eigenvalues are orthogonal, so the probability amplitude is

$$\begin{aligned} \langle \phi | e^{-iHt} | \psi \rangle &= (b_1^* \langle j_1, m_1 | + b_2^* \langle j_2, m_2 | + \dots) (a_1 | j_1, m_1 \rangle + a_2 | j_2, m_2 \rangle + \dots) \\ &= b_1^* a_1 M_{j_1} + b_2^* a_2 M_{j_2} + \dots \\ &\text{(states are normalized)} \end{aligned}$$

• Our examples using isospin

+ We can consider combinations of pions and nucleons and their isospin states (total)

$$\begin{aligned} \pi^+ + p &= | \frac{3}{2}, +\frac{3}{2} \rangle, & \pi^- + n &= | \frac{3}{2}, -\frac{3}{2} \rangle, \\ \pi^- + p &= (| \frac{3}{2}, -\frac{1}{2} \rangle - \sqrt{2} | \frac{1}{2}, -\frac{1}{2} \rangle) / \sqrt{3}, \\ \pi^0 + n &= (\sqrt{2} | \frac{3}{2}, -\frac{1}{2} \rangle + | \frac{1}{2}, -\frac{1}{2} \rangle) / \sqrt{3} \end{aligned}$$

+ It is impossible to have $\pi^+ + p \rightarrow \pi^0 + n$ or $\pi^- + p$ b/c the I_3 eigenvalue changes (also charge is not conserved)

+ However, the amplitudes

$$\begin{aligned} M(\pi^+ + p \rightarrow \pi^+ + p) &= \langle \frac{3}{2}, +\frac{3}{2} | M_{3/2} | \frac{3}{2}, +\frac{3}{2} \rangle = M_{3/2} \\ \text{and } M(\pi^- + n \rightarrow \pi^- + n) &= \langle \frac{3}{2}, -\frac{3}{2} | M_{3/2} | \frac{3}{2}, -\frac{3}{2} \rangle = M_{3/2} \\ &\text{are equal!} \end{aligned}$$

+ We can also find the other scattering amplitudes in a similar way. The same calculations work for decays or can work for angular momentum

+ There is a generalization for other symmetries like $SU(3)$