

Rigid Body Rotation

Math: Matrices + Tensors

- Index Notation (most examples in 3D, but more general)

• Suppose we choose a basis set of orthonormal vectors

$$\hat{i}, \hat{j}, \hat{k}$$

+ We remember we can write any vector \vec{a} as

$$\vec{a} = a_1 \hat{i} + a_2 \hat{j} + a_3 \hat{k}$$

+ The a_i numbers are the components of \vec{a} in this basis

+ We can assemble the components in a column matrix $\vec{a} \rightarrow \begin{bmatrix} a_1 \\ a_2 \\ a_3 \end{bmatrix}$

+ The dot product is $\vec{a} \cdot \vec{b} = a_1 b_1 + a_2 b_2 + a_3 b_3$
 $= \sum_i a_i b_i \equiv a_i b_i$

The last form is Einstein summation convention which means you sum over a repeated (paired) index

+ We can define the Levi-Civita symbol such that

$$\epsilon_{123} = 1, \quad \epsilon_{ijk} = -\epsilon_{jik} = -\epsilon_{ikj} = -\epsilon_{kji}$$

(minus sign every time you swap an index)

Then the cross product is

$$(\vec{a} \times \vec{b})_i = \sum_j \epsilon_{ijk} a_j b_k$$

+ Can also write the dot product as a matrix multiplication

$$[a_1 \ a_2 \ a_3] \begin{bmatrix} b_1 \\ b_2 \\ b_3 \end{bmatrix} = [\vec{a}]^T [\vec{b}]$$

where T = transpose = flips the column into a row.

• Matrices + vectors

+ Components according to rotated unit vectors are related by a rotation matrix

$$\vec{a} \rightarrow \begin{bmatrix} a'_1 \\ a'_2 \\ a'_3 \end{bmatrix} = R \begin{bmatrix} a_1 \\ a_2 \\ a_3 \end{bmatrix}$$

- + Other linear transformations are also matrix multiplication
- + In index notation, a matrix has an index indicating row position + one for column position

$$M_{ij} \quad \text{or} \quad M = \begin{matrix} \text{row} \\ \#1 \\ \#2 \\ \#3 \end{matrix} \begin{bmatrix} M_{11} & M_{12} & M_{13} \\ M_{21} & M_{22} & M_{23} \\ M_{31} & M_{32} & M_{33} \end{bmatrix} \begin{matrix} \text{column} \\ \#1 \\ \#2 \\ \#3 \end{matrix}$$

- + Matrix multiplication is a sum of column + row numbers

$$(AB)_{ij} = \sum_k A_{ik} B_{kj} \quad ; \quad (M\vec{a})_i = \sum_j M_{ij} a_j$$

- + The matrix transpose reverses row + column

$$(M^T)_{ij} = M_{ji} \quad \text{If } M^T = M, \text{ } M \text{ is symmetric}$$

So we can see $(AB)^T = B^T A^T$

- + Rotating axes should leave dot products unchanged

$$(\vec{a} \cdot \vec{b}) \rightarrow [a_1 \ a_2 \ a_3] \begin{bmatrix} b_1 \\ b_2 \\ b_3 \end{bmatrix} = [a_1 \ a_2 \ a_3] R^T R \begin{bmatrix} b_1 \\ b_2 \\ b_3 \end{bmatrix}$$

$$\Rightarrow R^T R = \mathbf{I} \quad \text{Means rotation matrices are orthogonal}$$

- Eigenvalues + eigenvectors

- What happens when we multiply a ^{square} matrix on a vector?

+ Normally, the components get mixed up
 + For special vectors \vec{a} given a matrix M ,

$$M\vec{a} = \lambda\vec{a} \quad \text{where } \lambda \text{ is a particular number.}$$

In other words, M just rescales \vec{a} .

- Definition: Given a ^{square} matrix M , a nonzero vector \vec{a}

is an eigenvector of M if $M\vec{a} = \lambda\vec{a}$
for some number λ called the eigenvalue

+ Typically, an $n \times n$ matrix has n distinct eigenvectors,
each with their own eigenvalue.

+ If M is (real) symmetric, eigenvectors can
be made orthonormal, and eigenvalues are real.

+ So the eigenvectors of a symmetric matrix make up
a basis of unit vectors

• How do we find eigenvalues and eigenvectors?

+ We can rewrite the basic equation as $(M - \lambda I)\vec{a} = 0$

+ $(M - \lambda I)$ is not invertible when λ is an eigenvalue
 $\Rightarrow \det(M - \lambda I) = 0$

+ This is a polynomial of order n for the eigenvalue λ
called the characteristic equation. The n roots $\lambda_1, \dots, \lambda_n$
are the n eigenvalues. They can repeat.

+ Then $M\vec{a} = \lambda_i\vec{a}$ is a system of linear eqns
for the components of the eigenvector \vec{a} with
eigenvalue λ_i . Then do this for each eigenvalue

+ Note: You can't determine the normalization of
an eigenvector from $M\vec{a} = \lambda\vec{a}$ (multiplying both
sides by a scalar is the same as rescaling \vec{a})
So add a normalization condition

• Some notes on complex matrices + vectors

+ The dot product for complex vectors is
 $\vec{a} \cdot \vec{b} = a_i^* b_i$ (often written (\vec{a}, \vec{b}))
b/c order matters.

+ The Hermitian adjoint $A^\dagger = (A^T)^*$, where $*$ is the complex conjugate of each element

+ A unitary matrix has $U^\dagger = U^{-1}$

A Hermitian matrix has $H^\dagger = H$

+ A Hermitian matrix has all real eigenvalues, and the eigenvectors can be chosen orthogonal

- Tensors

• Consider two vectors multiplied side by side but not dotted, i.e., $\vec{a} \vec{b} = \vec{T}$

+ This has 2 indices $T_{ij} = a_i b_j$ for components w.r.t. a set of unit vectors

+ It acts on a vector by dotting

$$(\vec{T} \vec{c}) = \vec{a} (\vec{b} \cdot \vec{c}) \quad \text{or} \quad \sum_j a_i b_j c_j$$

+ In general!

• A physical object (2nd) converts a vector to another vector (linearly) is a (rank-2) tensor

+ These can be "tensor products" of vectors or not

+ One different example is the identity 1 w/ components δ_{ij}

+ Given a set of axes, the components of tensors form a matrix, which acts on vectors by matrix multiplication on components $\vec{T} \vec{a} \rightarrow [T][a]$

+ So tensors have eigenvectors + eigenvalues

• How do tensor components change when axes are rotated?

+ Think about our vector product example.

For $T_{ij} = a_i b_j$,

$$T'_{ij} = \sum_{k,l} (R_{ik} a_k) (R_{jl} b_l) = R_{ik} R_{jl} T_{kl}$$

+ So the components of a tensor transform by rotating each index like a vector.

+ If we think of the components of a tensor as a matrix, we can re-organize the above

$$T'_{ij} = \sum_{k,l} R_{ik} T_{kl} (R^T)_{lj}$$

or $\vec{T}' \rightarrow [T'] = R[T]R^T$

+ What does this mean for components of $(\vec{T} \cdot \vec{a})$?

$$[T'] [a'] = R [T] R^T R [a] = R [T] [a]$$

In other words $(\vec{T} \cdot \vec{a})$ is a vector, so its components rotate like a vector's components