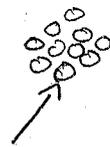


Scattering + Cross Sections

The Mean Free Path

- Suppose a particle moves through a region of other particles that present cross-sectional areas



+ If you track the particle's path through the material, how many other particles does it hit?

Reverse it: moving particle has cross-section σ , hits anything in cylinder around path



+ If the targets have number density n , # of collisions in length x is $= n\sigma x \Rightarrow$ average distance per collision is $\lambda = 1/n\sigma$. This is how far you expect a particle to go.

- Suppose there is a flux $f(x)$ of particles hitting the material. How does $f(x)$ change? Note: $f(x) = \text{particle/area in beam/time}$.

+ Each collision removes a particle. But we just saw that the total collisions/time/area is $f(x)(n\sigma dx)$

+ Therefore, $f(x+dx) - f(x) = \frac{df}{dx} dx = -f n\sigma dx \Rightarrow \frac{df}{dx} = -f/\lambda$
 $\Rightarrow f(x) = f_0 e^{-x/\lambda}$

- Turning this around, to the 1st picture again, consider a target and a flux f of incoming particles $\Rightarrow \Rightarrow \Rightarrow \circ$

+ The cross-section is really a property of the scattering, but treat it as belonging to the target

+ We have defined the target cross-section σ the ratio of the rate of collisions to the incoming flux
 $\# \text{ collisions/time} = f\sigma$

- Differential Cross Sections

• Cross Section is collisions / incoming flux. Want more specificity

+ If the initial velocity is along z axis w/ target at origin, the outgoing velocity is described by spherical polar angles θ, ϕ



+ We want # of scatterings with final velocity directed into a solid angle $d\Omega = \sin\theta d\theta d\phi$ at θ, ϕ

Recall: solid angle = area of region on spherical surface / radius².
4 π solid angles on whole sphere, measured in steradians

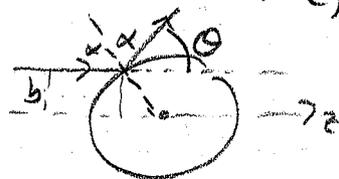
+ This # of scatterings $\propto d\Omega$ b/c also = $\int d\sigma$.

We define the differential cross section $d\sigma/d\Omega$ as # scatterings per incoming flux per solid angle. The total cross section is $\sigma = \int d\Omega \left(\frac{d\sigma}{d\Omega} \right)$.

• Ex Hard-sphere Scattering

+ Suppose the target is a sphere of radius R . The total cross section is $\sigma = \pi R^2$ b/c the flux sees a circular profile (assuming contact force)

+ Let's determine the scattering angle θ for particle incoming w/ impact parameter b .



For an elastic collision, the particle "reflects" around the radius. Therefore, $\theta = \pi - 2\alpha$, $\cos\alpha = b/R$.

$$\Rightarrow b = R \cos(\theta/2)$$

+ The flux that eventually scatters into a solid angle $d\Omega$ goes through area $d\sigma = |db| (b d\phi)$ where $db = -\frac{1}{2} R \sin(\theta/2) d\theta$ (This is negative b/c larger impact parameters scatter less.)

+ Substituting, $d\sigma = \frac{1}{4} R^2 \sin\theta d\theta d\phi \Rightarrow d\sigma/d\Omega = R^2/4$

This is isotropic and agrees with the expected total cross section.

◦ Ex Rutherford scattering

+ This is repulsive Coulomb potential scattering. Rutherford fired α particles (He nuclei) at gold foil. Electrons are too light to affect the α particle much, so $\frac{d\sigma}{d\Omega}(\theta)$ gives information on potential from nucleus. If nuclear charge fills the atom, potential is not Coulomb for small impact parameter

+ From our discussion of hyperbolic orbits, the impact parameter and scattering angle satisfy $b = \frac{k}{mv^2} \cot(\theta/2)$ with $k = \frac{q_1 q_2}{4\pi\epsilon_0}$ for the repulsive Coulomb potential

+ Then $d\sigma = b |db| d\phi$, and $db = \frac{-k}{2mv^2} \csc^2(\theta/2) d\theta$

With a little algebra, $\frac{d\sigma}{d\Omega} = \frac{1}{4} \left(\frac{k}{mv^2}\right)^2 \frac{1}{\sin^4(\theta/2)}$.

This is known as the Rutherford scattering cross section

+ The α particles can scatter to large angles, meaning there is a strong force at short distances \Rightarrow nucleus is very small.

+ The total cross section is divergent. That's because the Coulomb force is infinite-range. There is always a tiny bit of scattering.