

## PHYS-4601 Homework 2 Due 20 Sept 2018

This homework is due in the dropbox outside 2L26 by 10:59PM on the due date. You may alternately email a PDF (typed or black-and-white scanned) or give a hardcopy to Dr. Frey.

### 1. Diagonalization *Based on Griffiths A.26*

Consider a three-dimensional Hilbert space with orthonormal basis  $|e_i\rangle$ ,  $i = 1, 2, 3$ . The operator  $A$  takes the matrix representation

$$A = \sum_{i,j} |e_i\rangle\langle e_i| A |e_j\rangle\langle e_j| \simeq \begin{bmatrix} 2 & i & 1 \\ -i & 2 & i \\ 1 & -i & 2 \end{bmatrix}. \quad (1)$$

You should be able to check yourself that  $A$  is Hermitian.

- Find the eigenvalues  $a_i$  and corresponding eigenstates  $|a_i\rangle$  ( $A|a_i\rangle = a_i|a_i\rangle$ ) written in terms of their components  $\langle e_j|a_i\rangle$ . Choose the eigenstates to form an orthonormal eigenbasis; that is, choose any ambiguities such that  $\langle a_i|a_j\rangle = \delta_{ij}$ .
- As we will state in class,  $A$  can be written in the form

$$A = \sum_i a_i |a_i\rangle\langle a_i|, \quad (2)$$

where  $a_i$  are the eigenvalues and  $|a_i\rangle$  are the eigenvectors of  $A$ . Verify that formula (2) gives the same operator as (1) when you plug in your answer to part (a) for the eigenvalues and eigenvectors.

### 2. The Momentum Operator

The usual definition of the momentum operator  $p$  is that it acts by  $p \simeq -i\hbar d/dx$  in the position basis (precisely,  $\langle x|p|\psi\rangle = -i\hbar d\psi/dx$  for any state  $|\psi\rangle$  with wavefunction  $\langle x|\psi\rangle = \psi(x)$ ). In this problem, we will explore some properties of this operator and its eigenfunctions; in the future, we will see why it makes sense to call it momentum. For simplicity, we work in one dimension with  $-\infty < x < \infty$ .

- Let the state  $|p\rangle$  be an eigenstate of the momentum operator with real eigenvalue  $p$  (ie,  $p \cdot |p\rangle = p|p\rangle$ ). Show that  $|p\rangle$  has wavefunction

$$\langle x|p\rangle = \frac{1}{\sqrt{2\pi\hbar}} e^{ipx/\hbar} \quad (3)$$

(you may assume the normalization constant is given).

- Show that  $\langle p'|p\rangle = \delta(p - p')$ . *Hint:* You may find the formula

$$\delta(z) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dk e^{ikz} \quad (4)$$

helpful.

- Show that the wavefunction  $\psi(x) = \langle x|\psi\rangle$  and “momentum-space wavefunction”  $\tilde{\psi}(p) = \langle p|\psi\rangle$  for any vector  $|\psi\rangle$  are Fourier transforms, as defined in Griffiths equation [2.102] (up

to factors of  $\hbar$ ). To work this out precisely, it will be helpful for you to rescale  $x$  and  $p$  to remove explicit powers of  $\hbar$ .

What we have seen so far is that  $p \simeq -i\hbar d/dx$  has complex exponentials for eigenfunctions, and that this implies that  $\langle x|\psi\rangle$  and  $\langle p|\psi\rangle$  are Fourier transforms for any state  $|\psi\rangle$ . Now we want to prove the reverse: assuming (3), we will show that  $p \simeq -i\hbar d/dx$  in the position eigenbasis.

- (d) For momentum to be observable, it must be Hermitian. Assuming  $p$  is a Hermitian operator, show  $\langle p|p|\psi\rangle = p\tilde{\psi}(p)$  for any state  $|\psi\rangle$ . In other words, you are showing that  $p \cdot \tilde{\psi}(p) = p\tilde{\psi}(p)$  on momentum-space wavefunctions (like  $x \cdot \psi(x) = x\psi(x)$  on normal wavefunctions).
- (e) Assuming (3), we know that the wavefunction and momentum-space wavefunction are Fourier transforms of each other. Use this fact to show that

$$\langle x|p|\psi\rangle = -i\hbar \frac{d\psi}{dx}(x) . \quad (5)$$

What this means is that defining  $p \simeq -i\hbar d/dx$  is equivalent to defining the state  $|p\rangle$  by (3) — you can derive one statement from the other.

### 3. The Last Eigenvector

A system with a three-dimensional Hilbert space has Hamiltonian represented by the matrix

$$H \simeq \frac{E_0}{3} \begin{bmatrix} 1/2 & -1/2 & 2i \\ -1/2 & 1/2 & -2i \\ -2i & 2i & -1 \end{bmatrix} . \quad (6)$$

Two of the eigenstates are represented by the column vectors

$$|1\rangle \simeq \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ 1 \\ 0 \end{bmatrix} \quad \text{and} \quad |2\rangle \simeq \frac{1}{\sqrt{3}} \begin{bmatrix} i \\ -i \\ 1 \end{bmatrix} . \quad (7)$$

- (a) Find the third eigenstate  $|3\rangle$  as a column vector, including proper normalization.
- (b) Find all the eigenvalues of  $H$  [EXTRA CREDIT: and write  $H$  as a matrix in the  $\{|1\rangle, |2\rangle, |3\rangle\}$  basis.]
- (c) Write the projection operator  $|3\rangle\langle 3|$  as a matrix in the original basis [EXTRA CREDIT: and the  $\{|1\rangle, |2\rangle, |3\rangle\}$  basis.]