PHYS-4601 Homework 17 Due 2 Mar 2017

This homework is due in the dropbox outside 2L26 by 10:59PM on the due date. You may alternately email a PDF (typed or black-and-white scanned) or give a hardcopy to Dr. Frey.

Useful information for this assignment:

From our discussion of the free electron gas, the number of states of a free particle with wavevector magnitude between k and k + dk in a volume V is

$$\frac{g}{8}\frac{4\pi k^2 dk}{\pi^3/V} = \frac{gV}{2\pi^2}k^2 dk , \qquad (1)$$

where g is the number of spin or polarization states. This is one octant of a spherical shell of radius k and thickness dk with one state per π^3/V . This number of states is true for any type of free particle, relativistic or not. You may set Boltzmann's constant $k_B = 1$ for this assignment.

1. An Integral, Gamma, and Zeta Functions

The gamma function and Riemann zeta function are defined as

$$\Gamma(z) = \int_0^\infty dx \, x^{z-1} e^{-x} \text{ and } \zeta(s) = \sum_{n=1}^\infty n^{-s} .$$
 (2)

For use on the rest of the assignment, prove that

$$\int_0^\infty dx \, \frac{x^{s-1}}{e^x - 1} = \Gamma(s)\zeta(s) \ . \tag{3}$$

In the remaining problems, use equation (3) when appropriate and note that $\Gamma(n+1) = n!$ for n a non-negative integer. *Hint:* expand the fraction as a geometric series $\sum_{n=0}^{\infty} x^n = 1/(1-x)$.

2. Bose-Einstein Condensation from Griffiths 5.29

Consider the Bose-Einstein distribution for spin-0 bosons. As in equation (1), we can usually approximate the sum over states as an integral over wavevectors so that the number density and energy density are

$$n = \frac{1}{2\pi^2} \int_0^\infty dk \, k^2 f(x) \,, \quad \rho = \frac{\hbar^2}{4m\pi^2} \int_0^\infty dk \, k^4 f(x) \,, \quad x = \frac{1}{T} \left(\frac{\hbar^2 k^2}{2m} - \mu \right) \,, \tag{4}$$

where f is the Bose-Einstein distribution $f(x) = 1/(e^x - 1)$. Note that a spin-0 boson has g = 1 spin state.

If there is a fixed number density n of bosons, the chemical potential μ is a function of T. $\mu = 0$ at a fixed critical temperature T_c ; $\mu < 0$ for higher temperatures, and the integral approximation above fails for lower temperatures because all the bosons go into the ground state (this is *Bose-Einstein condensation*). Find T_c as a function of n. You may write your answer in terms of gamma and zeta functions.

3. Proton Chemical Potential in Cosmology

For much of the history of the universe, protons are nonrelativistic and in thermal equilibrium with photons. During that time period, the number densities of protons n_B and photons n_{γ} are determined by the Fermi-Dirac and Bose-Einstein distributions with the same temperature. (a) Photons with g = 2 polarization states have single particle energies given by $\epsilon = c\hbar k$. Assuming zero chemical potential, show that the number density of photons at temperature T is

$$n_{\gamma} = \frac{2\zeta(3)}{\pi^2} \left(\frac{T}{\hbar c}\right)^3 \ . \tag{5}$$

Hint: Remember, photons obey Bose-Einstein statistics.

- (b) Now consider nonrelativistic protons. Including the rest mass energy, a single particle state has energy $\epsilon = mc^2 + \hbar^2 k^2/2m$. Assume that $0 < \mu < mc^2$ and find the number density n_p of protons when $T \ll (mc^2 \mu)$. *Hint:* Argue that the 1 in the denominator of the distribution function is negligible in this limit.
- (c) The ratio $\eta = n_p/n_\gamma$ is known as the *baryon-to-photon* ratio in cosmology (ignoring the small proportion of elements besides hydrogen), and it is constant over the period of the universe we consider with a value of $\eta \approx 10^{-10}$. Find the chemical potential μ of the protons in terms of η , the temperature T, the proton mass m, and physical constants. You may ignore constants of order unity.