

PHYS-4601 Homework 14 Due 30 Jan 2014

This homework is due in the dropbox outside 2L26 by 11:59PM on the due date. If you wish to turn it in ahead of time, you may email a PDF or give a hardcopy to Dr. Frey.

1. Bloch's Theorem

In class, we stated Bloch's theorem as saying that any stationary state wavefunction of a periodic potential $V(x+a) = V(x)$ can be written to satisfy the condition

$$\psi(x+a) = e^{iKa}\psi(x) \quad (K \text{ real}) . \quad (1)$$

Alternately, Bloch's theorem can state that any stationary state wavefunction of the same potential can be written as

$$\psi(x) = e^{iKx}u(x) \text{ where } u(x+a) = u(x) . \quad (2)$$

- Show that these two formulations are equivalent (that is, show that if ψ satisfies (1) then it satisfies (2) and vice-versa).
- If you write $u(x) = \sum_q c_q e^{iqx}$ as a Fourier series, what are the allowed values of q ? Therefore, what are the allowed values of the momentum for a wavefunction with a given K ?

2. Specific Heat of the Free-Electron (Fermi) Gas

In this problem, you'll explore the thermal properties of the electrons in a simple metal. Recall that the total number of electrons and total energy of the gas (on average) are

$$N = \sum_{\text{states}} f\left(\frac{\epsilon - \mu}{T}\right) , \quad U = \sum_{\text{states}} \epsilon f\left(\frac{\epsilon - \mu}{T}\right) , \quad f(x) = \frac{1}{e^x + 1} . \quad (3)$$

Here, T is the temperature, and μ is the chemical potential. For free electrons in volume V , the energy of each state is $\epsilon = \hbar^2 k^2 / 2m$ in terms of the wavevector magnitude k , and there are $(V k^2 / \pi^2) dk$ states at that energy. It is a good approximation to replace the sum over states with integrals to define the number and energy densities

$$n = \frac{1}{\pi^2} \int_0^\infty dk k^2 f(x) , \quad \rho = \frac{\hbar^2}{2m\pi^2} \int_0^\infty dk k^4 f(x) , \quad x = \frac{1}{T} \left(\frac{\hbar^2 k^2}{2m} - \mu \right) . \quad (4)$$

In this problem, you will find the specific heat of the free electron gas at low temperatures, which is defined as

$$c_V = \frac{d\rho}{dT} \text{ with } \frac{dn}{dT} = 0 . \quad (5)$$

Work only to lowest order in the temperature.

- Argue that you can approximate

$$\int_0^\infty dk k^2 g(x) = \frac{T}{2\mu} \left(\frac{2m\mu}{\hbar^2} \right)^{3/2} \int_{-\infty}^\infty dx g(x) \quad (6)$$

for $x = (\epsilon - \mu)/T$ and any function $g(x)$ as $T \rightarrow 0$. Be careful with limits of integration.

- Use conservation of particle density $dn/dT = 0$ to show that $d\mu/dT \rightarrow 0$ as $T \rightarrow 0$. Hint: Use your previous results and also the fact that df/dx is an even function of x .

(c) Show that

$$c_V = \frac{mT}{\hbar^3 \pi^2} \sqrt{2m\mu} \int_{-\infty}^{\infty} dx x^2 \frac{df}{dx}, \quad (7)$$

where the x integral is a pure number. To lowest order, you may set $d\mu/dT \rightarrow 0$ as follows from part (b).

(d) Finally, at $T = 0$, $\mu = E_F$, the Fermi energy, as discussed in class. Therefore argue that $c_V \propto Tn^{1/3}$.