PHYS-4601 Homework 20 Due 28 Mar 2013

This homework is due in the dropbox outside 2L26 by 11:59PM on the due date. If you wish to turn it in ahead of time, you may email a PDF or give a hardcopy to Dr. Frey.

1. **The Path Integral and Ehrenfest's Theorem**

   As we've seen in the lecture notes, the probability amplitude of a transition from one given state at time t_i to another state at time t_f, as well as the expectation value of any variable $\mathcal{O}$ (called an operator in Hamiltonian quantum mechanics), can be given by the path integrals

   $$\int \mathcal{D}x \, e^{iS/\hbar} \; , \quad \langle \mathcal{O} \rangle = \int \mathcal{D}x \, \mathcal{O} e^{iS/\hbar} \; , \tag{1}$$

   where S is the action of the system, given appropriate boundary conditions and normalization (which we intentionally leave vague). There is one other property of functional integrals that we haven't discussed yet:

   $$\int \mathcal{D}x \, \frac{\delta}{\delta x(t)} \, (\text{anything}) = 0 \; , \tag{2}$$

   where $\delta/\delta x(t)$ is the functional derivative. This basically says that the integral of a total derivative is zero (we throw away boundary terms on configuration space; heuristically, we are integrating functions that go to zero as $x \to \pm\infty$).

   In this problem, consider a system with Lagrangian $L = (m/2)\dot{x}^2 - V(x)$ and action $S = \int dt \, L$.

   (a) Use (2) to show that the expectation value of x follows the classical Euler-Lagrange equation

   $$\left\langle \frac{\delta S}{\delta x(t)} \right\rangle = 0 \text{ or equivalently } m\frac{d^2}{dt^2}\langle x \rangle = -\left\langle \frac{dV}{dx}(x) \right\rangle \; . \tag{3}$$

   (b) Derive the Hamiltonian associated with the given Lagrangian. Then go back to our discussion of Ehrenfest's theorem from the fall (you can also find this in the Griffiths textbook). Show that the first order differential equations that determine $d\langle x \rangle/dt$ and $d\langle p \rangle/dt$ imply equation (3).

   (c) In this part, set the potential $V = 0$. By the method of your choice (*hint*: think about part (a)), show that the expectation value $\langle x(t)x(t') \rangle$ satisfies the differential equation

   $$\left(m\frac{d^2}{dt^2} \right) \langle x(t)x(t') \rangle = -i\hbar\delta(t - t') \; . \tag{4}$$

2. **The Double-Slit Experiment and Aharonov-Bohm Effect**

   This problem will consider the classic double-slit experiment from the path integral point of view along with a version with a twist.

   (a) In the double-slit experiment illustrated below, two paths dominate the path integral for particles emitted at A and measured at D: from A to B to D with constant speed and

from A to C to D at constant speed (the figure is not to scale).
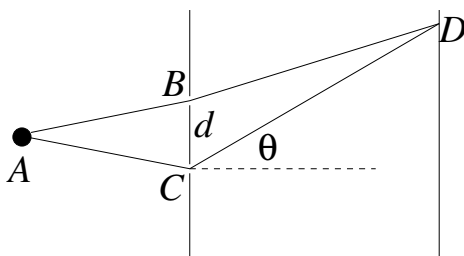
That means the amplitude for measurement of a particle at D is

$$\mathcal{M} = \int \mathcal{D}x \, e^{iS/\hbar} \approx e^{iS_{A\to B\to D}/\hbar} + e^{iS_{A\to C\to D}/\hbar} \; . \tag{5}$$

Assume that the two slits at B and C are separated by a distance d which is very small compared to the distance between the two screens, so that segments BD and CD make the same angle θ with respect to the horizontal, and ignore the size of the slits. Show that the probability of observing the particle at D is proportional to $\cos^2(\pi d \sin\theta/\lambda)$, where $\lambda = 2\pi\hbar/p$ is the de Broglie wavelength. *Hint:* The length of each path ABD or ACD is much longer than the difference in path lengths, so you will want to do an expansion in the difference. Remember that both paths must last the same length of time, and use the total path length to determine the momentum.

(b) Next, assume the particle has charge q but that the scalar potential is zero. Recall that the Lagrangian for a charged particle in a vector potential is $L = (m/2)\dot{\vec{x}}^2 + q\dot{\vec{x}} \cdot \vec{A}$ and show that the probability for measuring the particle at D becomes

$$P \propto \cos^2\left[\frac{\pi d}{\lambda}\sin\theta + \frac{q}{2}\oint \vec{A}\cdot d\vec{l}\right] , \tag{6}$$

where the integral is from $A \to C \to D \to B \to A$.

(c) Now consider a solenoid added to the experiment (the blue circle in the figure below) with a magnetic field pointed out of the paper. The particle cannot penetrate the solenoid.
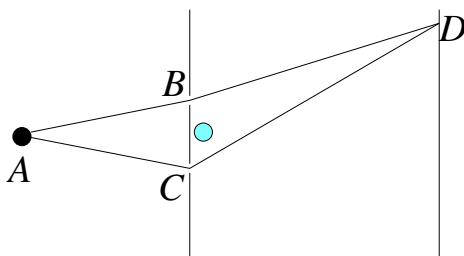
Now use Stokes's Theorem (*aka* the curl theorem; see Griffiths's EM Theory textbook if necessary) to show that

$$P \propto \cos^2\left[\frac{\pi d}{\lambda}\sin\theta + \frac{q}{2}\Phi_B\right] , \tag{7}$$

where Φ_B is the flux of the magnetic field through the solenoid. Amazingly enough, this shows that the magnetic field influences the interference pattern of the particle, even though the particle (or its wavefunction) cannot enter a region with a nonzero magnetic field! This is called the *Aharonov-Bohm effect*.