## PHYS-4601 Homework 17 Due 8 Mar 2012

This homework is due in the dropbox outside 2L26 by 11:59PM on the due date. If you wish to turn it in ahead of time, you may email a PDF or give a hardcopy to Dr. Frey.

## 1. More Notes on Time-Dependent Perturbation Theory

In this problem, consider a Hamiltonian  $H = H_0 + H_1(t)$ , where we know the eigenstates  $|\psi_n^0\rangle$  and eigenvalues  $E_n^0$  of  $H_0$  and where  $H_1(t)$  is a small time-dependent contribution to the Hamiltonian. We write the full time-dependent state as

$$|\Psi(t)\rangle = \sum_{n} c_n(t)e^{-iE_n^0 t/\hbar} |\psi_n^0\rangle . \tag{1}$$

- (a) In class, we derived the formula for  $c_n(t)$  to first order in perturbation theory, assuming that  $H_1(t)$  is small. Prove that those  $c_n(t)$  satisfy the normalization condition  $\sum_n |c_n|^2 = 1$  to first order in  $H_1$  at all times. (See the lecture notes for the appropriate formula; equation [9.17] in Griffiths is *not* general enough for this problem.)
- (b) Suppose there are just two  $H_0$  eigenstates  $|1\rangle = |\psi_1^0\rangle$  and  $|2\rangle = |\psi_2^0\rangle$  with perturbation Hamiltonian

$$\langle 1|H_1|1\rangle = \langle 2|H_1|2\rangle = 0$$
,  $\langle 1|H_1|2\rangle = \langle 2|H_1|1\rangle = \begin{cases} 0 & t < 0 \text{ or } t > T \\ V & 0 \le t \le T \end{cases}$ . (2)

If the initial state is  $|\Psi(t=0)\rangle = |1\rangle$ , find the probability that a measurement finds the system in state  $|2\rangle$  at time t=T to first order in  $H_1$ . Hint: You may directly do the integration or use the limit of the sinusoidal perturbation discussed in class.

## 2. NMR aka MRI easier version of Griffiths 9.20

Consider a spin-1/2 particle (for example, a proton) with gyromagnetic ratio  $\gamma$  in the presence of a magnetic field

$$\vec{B} = B_0 \hat{z} + B_1 \cos(\omega t) \hat{x} - B_1 \sin(\omega t) \hat{y} , \quad B_1 \ll B_0$$
(3)

at its fixed position. This is how an NMR (MRI) machine works; a proton sitting in a large static magnetic field is exposed to a small radio-frequency magnetic field. This problem is exactly solvable (as in the Griffiths problem), but you are to use first-order time-dependent perturbation theory.

(a) The Hamiltonian is given by the usual interaction between a magnetic moment and magnetic field,  $H = -\gamma \vec{B} \cdot \vec{S}$ . Show that the Hamiltonian can be written as  $H = H_0 + H_1(t)$ , where  $H_0$  has the same eigenstates as  $S_z$  with eigenvalues  $\pm \gamma B_0 \hbar/2$  and

$$H_1 = -\frac{\gamma B_1 \hbar}{2} \begin{bmatrix} 0 & e^{i\omega t} \\ e^{-i\omega t} & 0 \end{bmatrix} . \tag{4}$$

(b) At first order in the small field  $B_1$ , find the probability that a particle which is spin-up at t = 0 is measured to be spin-down at some later time t. For a given frequency  $\omega$ , at what times is this probability maximized (that is, if we want to flip as many proton spins as possible, how long should we leave the radio-frequency pulse on)?

## 3. Fermi's Golden Rule

Consider a sinusoidal perturbation Hamiltonian  $H_1 = Ve^{-i\omega t} + V^{\dagger}e^{+i\omega t}$ . In the class notes, we found the probability for a transition from state  $|1\rangle$  to  $|2\rangle$  as a function of time and frequency  $\omega$ . In the following, define  $\hbar\omega_0 = E_2 - E_1$ , the difference of the energy eigenvalues of the unperturbed Hamiltonian  $H_0$ . We will investigate the transition probability near  $\omega = \omega_0$  at large t (at least as long as the probability stays small).

- (a) At a fixed (and large) time, the probability is peaked at  $\omega = \omega_0$ . Using L'Hospital's rule or just a power series expansion, find the peak transition probability as a function of time.
- (b) Find the values of  $\omega$  where the probability first vanishes on either side of  $\omega = \omega_0$ . The difference in these two values tells us the width of the peak.
- (c) For large enough times, approximate the transition probability as a rectangle with the peak value from part (a) and width given by half the difference in part (b). Integrate this approximate probability function and argue that

$$P \to \frac{2\pi |V_{21}|^2}{\hbar^2} t\delta(\omega_0 - \omega) \tag{5}$$

as  $t \to \infty$ .

This problem shows two things: first, transitions occur only to states at energies related by the perturbation frequency and, second, that there is a constant transition rate (probability per unit time) to the appropriate states. The relationship (5) is known as *Fermi's Golden Rule*.